Solute Dispersion in Porous Media

Beyond Percolation Scaling

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What is dispersion?

• Migration apart of particles under influence of fluid flow and molecular diffusion

Where is it relevant?

"underpins processes as diverse as cellular mitosis, blood perfusion in the brain, chromatography, filtration, secondary oil recovery, ground water remediation, catalysis, and the behavior of packed bed reactors." degradation of building materials tissue physiology, migration and epidemiology.

Most common approach

- The advection-dispersion equation (ADE)
- It models solute transport on a continuum as due to diffusion like processes (D_h) and being carried along by the flow (advection) *v*.

$$\frac{\partial C}{\partial t} = D_h \nabla^2 C - \mathbf{v} \cdot \nabla C$$

Are there problems with the ADE?

- $D_{\rm h} >> D_{\rm m}$ has nothing to do with molecular diffusion, $D_{\rm m}$.
- $D_{\rm h}$ depends in a non-trivial way on velocity.
- D_h increases with time (or transport distance) in a way that cannot be predicted from the ADE.
- The distribution tails are seldom, if ever, Gaussian, which underestimates solute arrivals at both short and long times.
- All these features drop out of a proper accounting for effects of heterogeneity and connectivity on solute paths through the medium; first term is extraneous.
- The chief lacks of the ADE are admitted to be its treatment of heterogeneity and connectivity.

Anything good about ADE?

• Works at scale of single pore, where it is used to evaluate relative effects of diffusion and advection (Peclet number).

Why percolation theory?

- Previous approaches (e.g., the ADE) have failed to predict the observed behavior
- Percolation theory (PT) is a mathematics of pathways and connections in disordered materials
- Can treat heterogeneity with a wide range of flow resistance (critical path analysis, CPA), even far from percolation threshold
- It accounts for tortuosity of flow paths
- It can incorporate geological correlations

Organization

- Introduce concept of critical path analysis (CPA) to find the critical conductance g_c of a system.
 - 1. Apply CPA framework to cluster statistics of percolation theory to find spanning probability, W(g; x) that system of length x can be spanned by interconnected cluster of conductances of arbitrary smallest value, g.
 - 2. Apply critical scaling of percolation theory to find how the solute transit time, *t*, depends on the length of the system and the governing value, *g*.
 - 3. Arrival time distribution W(t;x) = g W(g;x) / (dt(g)/dg)
- Analogous procedure gives spatial distribution.
- Calculate moments by standard procedures.

First, some percolation theory concepts



On left a model porous medium.



In fractal medium we have "pore size distribution"

The biggest pores don't connect



Log pore size

The next-biggest still don't connect



Cluster with largest "volume"



Log pore size

Something important happened



Log pore size

In other words: Largest cluster reaches infinite size

How big is the bottleneck pore?



Infinite cluster (finite clusters have been deleted)

Bottleneck pore size Critical radius r_c



Log pore size

Bottleneck pore has conductivity g_c



Backbone

Bonds that actually conductDangling ends have been deleted

Bottleneck pore size Critical conductance g_c



Log pore size

$$W(r) = \frac{3-D}{r_m^{3-D}} r^{2-D} \quad \text{IMPLIES} \qquad W(g) \propto g^{\frac{-D}{3}}$$

$$\int_{g_c}^{g_m} W(g) dg = p_c$$

Analogously (for future information)

$$\int_{g}^{g_{m}} W(g) dg = p$$

We just got

• critical conductance, g_c

Percolation theory gives us

- cluster size statistics
- cluster topology

How can we work in PT using g and g_c ?

In percolation theory, quantities of interest are expressed in terms of p and p_c . Results of previous slides transform p and p_c to g and g_c . This transformation of variables allows representation of correlation length, cluster statistics, tortuosity, etc. in terms of a conductance value.

Cluster statistics

$$n_s \approx s^{-\tau} \exp\left\{-\left[(p-p_c)s^{\sigma}\right]^2\right\}$$

Use transformation and Let N ^{1/ σ v} = s. Use $\frac{\tau - 1}{\sigma v} = d$

$$n_{N} = \frac{1}{N^{d+1}} \exp\left\{-\left[\left(\frac{Nl}{L}\right)^{\frac{1}{\nu}} \left|1 - \left(\frac{g}{g_{c}}\right)^{\frac{3-D}{3}}\right|\right]^{2}\right\}$$

Result (after integrating over all clusters larger than system):

$$W(g;x) \propto \frac{1}{\beta} Ei \left[\alpha \left(\frac{x}{L} \right)^{\beta} \right]$$

$$Ei[z] = \int_{z}^{\infty} \frac{\exp[-y]}{y} dy$$

$$\alpha = \left| 1 - \left(\frac{g}{g_c} \right)^{\frac{3-D}{3}} \right|^2 \quad \text{and} \quad \beta = \frac{2}{\nu}$$

Result is not a power-law, exponential integral has logarithmic divergence at $g = g_c$



Calculate *t*(*g*), the time particles take to traverse a cluster characterized by *g*.

There are two contributions: one from the conductance distribution W(g), and one from topology via percolation theory. $t = \left(\frac{x}{L}\right)^{D_b} \frac{t_0}{3-D} \frac{1}{(1-V_c)^{\Lambda-v}} \left[\left(1 + \frac{V_c}{1-V_c}\right) \left(\frac{g_c}{g}\right)^{1-D/3} - 1 \right] \left[\frac{1}{\left[\left(\frac{g}{g_c}\right)^{1-D/3} - 1 \right]} \right]^{(D_b-1)v}$ $= \left(\frac{x}{L}\right)^{D_b} t_g$



Apply $W_{p}(t; x) = g W_{p}(g; x) / (dt(g)/dg)$





Note that we do not obtain power-law tail, but our results match such experimental results rather closely. Example given has slope -1.56, almost identical to what is seen in fracture flow experiments.

2-D flow, parameters from 2-D random percolation. Navier-Stokes simulations on 2-D structure at Percolation threshold



No prediction here, only a fit, but to "classic" Nielsen and Biggar experiment (cited by Cortis and Berkowitz)



Analogous procedure for spatial distribution at time t



Define longitudinal dispersion coefficient and dispersivity

$$D_{\rm I}(t) \equiv \frac{1}{2} \frac{d}{dt} \left[\frac{x^2}{2} - \frac{x^2}{2} \right] \approx \left[\frac{x^2}{2} - \frac{x^2}{2} \right] / t.$$

$$\alpha \equiv \left[< x^2 > - < x > 2 \right] / < x >.$$

Averages performed over distributions derived above: for Gaussian spreading, $D_{\rm I}(t)$ would be constant in time, making α constant in space.



























Length Dependence of Typical System Crossing Time

Power	Predicted	Reference	Material
1.67	1.64	Pfister and Griffiths ¹²⁷	Carbazole polymers
1.72	1.64	Pfister and Griffiths ¹²⁷	Carbazole polymers
2.1	1.9	Pfister ¹²⁸	As ₂ Se ₃
2	1.9	Tiedje ¹²⁹	a-Si:H
2.2	can't predict	Scher and Montroll ⁶⁵	not given
1.91	1.9	Pfister ¹³⁰	a-Se
1.86	1.9	Pfister ¹³⁰	a-Se
1.61	1.64	Bos et al ¹³¹	Polyvinyl carbazole
1.82	can't predict	Pfister and Scher ¹³²	not given

Summary

- Straightforward derivation in terms of known quantities from percolation theory leads to:
- Arrival time distribution consistent with simulations and experiments in soils and fracture flow
- Superlinear dependence of variance on time (at short times) in accord with experimental summaries and possibly predictive for Borden aquifer
- Dependence of dispersivity on length scale in accord with over 2400 experiments over 10 orders of magnitude of length scale
- Dependence of dispersivity on heterogeneity which is in accord with experiment for small heterogeneity, but which may or may not be for large heterogeneity
- Dependence of typical transport time on system length which is observed in semiconductors/polymers.

Implications

- No need to invoke multiple scales of heterogeneity to produce continuing rise of dispersivity with length scale (trouble for stochastic subsurface hydrology and NSF).
- Tendency for all experiments at scales of centimeters and above to have same REV (1m³) implies the relevance of experimental apparatus (human size on order of 2m).
- No role of diffusion means trouble both for ADE users and for proponents of matrix diffusion to explain solute transport retardation in fracture flow (nuclear regulatory issue).
- Where is the importance of geological complexity?